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# Well Injectivity Decline for Nonlinear Filtration of Injected Suspension: Semi-Analytical Model

Injectivity decline due to injection of water with particles is a widespread phenomenon in waterflood projects. It happens due to particle capture by rocks and consequent permeability decline and also due to external cake formation on the sandface. Since offshore production environments become ever more complex, particularly in deep water fields, the risk associated with injectivity impairment due to injection of seawater or re-injection of produced water may increase to the point that production by conventional waterflood may cease to be viable. Therefore, it is becoming increasingly important to predict injectivity evolution under such circumstances. The work develops a semi-analytical model for injectivity impairment during a suspension injection for the case of filtration and formation damage coefficients being linear functions of retained particle concentration. The model exhibits limited retained particle accumulation, while the traditional model with a constant filtration coefficient predicts unlimited growth of retained particle concentration. The developed model also predicts the well index stabilization after the decline period. [DOI: 10.1115/1.4002242]

Keywords: injectivity, formation damage, analytical model, well index, deep bed filtration

### **1** 1 Introduction

2 Injectivity decline is widely spread during waterflooding. One 3 of the main reasons for the phenomenon is the capture of solid and 4 liquid particles from injected aqueous suspension (seawater, pro-5 duced water, or any poor quality water) resulting in reduction of 6 permeability and of injection well index [1–4]. The reliable pre-7 diction of injectivity index decline is important for the planning 8 and design of well fracturing, acidizing, and other well stimula-9 tion techniques. It is also important for the design of water man-10 agement strategy for waterflood projects.

Injection well behavior prediction under rock clogging due to injection of the colloid/suspension of particles is based on the mathematical modeling of particle transport in porous media [5,6]. The main parameter determining kinetics of particle capture by a porous matrix is the filtration coefficient, which is the ratio between the particle retention rate and the module of the particle flux [6]. The filtration coefficient is a probability for particles to be captured during its transport over the unity distance. Dependency of the filtration coefficient on the retained particle concentration is called the filtration function. Its form highly affects the injectivity decline.

22 Different forms of filtration function have been presented in the 23 literature [6,7]. The constant filtration coefficient corresponds to 24 low retained concentrations, i.e., to the beginning of the well im-25 pairment process. The linear filtration function (the so called 26 blocking function) corresponds to Langmuir particle deposition, 27 where the retention rate is proportional to the number of vacant 28 sites in the porous space and the matrix surface. The blocking 29 filtration function is realized under intermediate retained concen-30 trations [2,8]. Different nonlinear shaped filtration functions have 31 been observed for highly clogged rocks [7,8]. Analytical models for deep bed filtration have been presented 32 for the case of constant filtration and formation damage coeffi-33 cients for both linear flow geometry (laboratory corefloods) and 34 axisymmetric flow (flow around the vertical well) [6,9,10]. The 35 models exhibit unlimited growth of retained particle accumulation, which is physically unrealistic. Nevertheless, either coreflood 37 or well data show the stabilized permeability with time. The analytical model for coreflood with blocking filtration function also 99 exhibits the stabilized injectivity index [6,11,12].

In the current work, a semi-analytical model for axisymmetric 41 flow around the vertical well is developed for a blocking filtration 42 function. The particle capture stops after reaching the maximum 43 (critical) value by the retained concentration in each point of the 44 drainage volume. The injectivity index stabilizes when the re- 45 tained concentration grows up to the critical value in a well neigh- 46 borhood.

The structure of the paper is as follows. First, we describe the **48** mathematical model for suspension filtration and rock clogging **49** under Langmuir particle retention. Then, we discuss a semi- **50** analytical model for axisymmetric flow with detailed derivations. **51** Afterward, the results of impedance prediction with a sensitivity **52** analysis are presented. **53** 

### 2 Mathematical Model for Deep Bed Filtration 54

During seawater injection, mainly solid particles penetrate into 55 the reservoir; their retention in rock results in permeability decline 56 and a consequent decrease in well injectivity index. Oily water 57 injection during produced water re-injection (PWRI) also results 58 in injectivity impairment. 59

Deep bed filtration with similar modeling challenges occurs 60 during drilling fluid invasion into formation with consequent per- 61 meability damage [13–15]. It also occurs during produced water 62 disposal in aquifers [16,17]. 63

It is assumed that all injected particles have the same filtration 64 properties, allowing us to introduce the overall suspended and 65 retained particle concentrations. Figure 1 shows the suspension 66

Contributed by the Petroleum Division of ASME for publication in the JOURNAL OF ENERGY RESOURCES TECHNOLOGY. Manuscript received January 1, 2009; final manuscript received June 7, 2009; published online xxxxx-xxxxx. Assoc. Editor: Desheng Zhou.

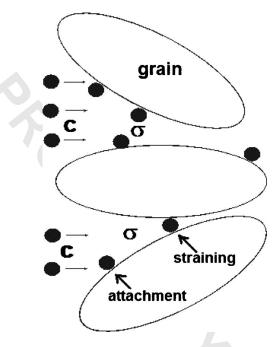


Fig. 1 Suspended and retained particle concentrations in porous space

67 concentration c of moving particles and the concentration  $\sigma$  of 68 particles attached to grain surfaces and captured in thin throats by 69 size exclusion.

70 Following Refs. [4,6,18,19], let us describe the mathematical 71 model for flow of suspensions in porous media in the axisymmet-72 ric geometry of vertical wells.

73 The conservation of suspended and retained particles in porous74 media is

(1)

(2)

(3)

 $\phi \frac{\partial c}{\partial t} + \frac{q}{2\pi rh} \frac{\partial c}{\partial r} = -\frac{\partial \sigma}{\partial t}$ 

76 It is assumed that water is incompressible, and particle retention 77 by rock does not change the total volume of the system "water 78 particles." It results in the conservation of the flux q.

 Classical filtration theory assumes that the retention rate is pro- portional to the particle advective flux cU, i.e., increase in either suspended concentration or velocity in number of times results in the same number of times increase in the number of particles, which are transported to vacancies in rocks per unit of time. The proportionality coefficient  $\lambda$  depends on  $\sigma$  and is called the filtra-tion function  $\lambda(\sigma)$ ,

$$\frac{\partial \sigma}{\partial t} = \lambda(\sigma) c \frac{q}{2\pi rh}$$

87 We assume that the capture rate is also proportional to the num-88 ber of vacancies (Langmuir's hypothesis). If A is a specific rock 89 surface and b is an "individual" area on the grain surface engaged 90 by one retained particle, the vacancy concentration is proportional 91 to the free vacant surface, which equals  $A-b\sigma$  (Fig. 1). 92 Finally, the retention rate is proportional to

 $(A - b\sigma)cU$ 

94 Introducing the proportionality coefficient yields the kinetic95 equation for retention rate,

 $\frac{\partial \sigma}{\partial t} = \lambda_0 \left( 1 - \frac{\sigma}{\sigma_m} \right) c \frac{q}{2\pi rh}$ 

96

93

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The comparison between Eqs. (2) and (3) shows that the filtra- **97** tion function in Eq. (3) is linear with respect to retention concen- **98** tration, **99** 

$$\Lambda(\sigma) = \begin{cases} \lambda_0 \left( 1 - \frac{\sigma}{\sigma_m} \right), & \sigma < \sigma_m \\ 0, & \sigma > \sigma_m \end{cases}$$
 100

If linear dependency in Eq. (3) for  $\lambda(\sigma)$  holds for the whole 101 interval of retention concentration variation, the maximum of re- 102 tention concentration  $\sigma_m$  corresponds to the case where the overall 103 vacant grain surface is filled by particles; i.e.,  $\sigma_m$  is proportional 104 to the initial number of vacancies. If  $\lambda(\sigma)$  is linear just for some 105 initial interval of  $\sigma$  and becomes nonlinear for larger values of 106 retained concentration,  $\sigma_m$  is just a coefficient in the equation for 107 a straight line without any specific physics meaning. 108

So, the retention rate is characterized by two constants—by the **109** initial filtration coefficient  $\lambda_0$  and by the maximum retention con- **110** centration  $\sigma_m$ . Parameters  $\lambda_0$  and  $\sigma_m$  depend on the salinity and **111** pH of the injected water, on the mineral composition of grains, on **112** particle size, on its wettability, on temperature, etc.; i.e., two con- **113** stants are determined by the rock and the injected fluid. **114** 

If the retention concentration is negligibly smaller than the ini- 115 tial concentration of vacancies, the retained particles do not affect 116 the vacancy concentration and do not change the retention condi- 117 tions, and the filtration function is constant. Formally, it corre- 118 sponds to an infinite value of the maximum retention concentra- 119 tion  $\sigma_m$ .

So called collective effects of particle interaction at high con- 121 centrations lead to the nonlinear filtration coefficient  $\lambda = \lambda(\sigma)$ . Its 122 form depends on brine salinity, pH, electric DLVO forces, etc. 123 [20-22]. 124 AQ:

Particle retention results in permeability decrease; i.e., perme- 125 ability is  $\sigma$  dependent: 126

$$\frac{q}{2\pi rh} = -\frac{k_0 k(\sigma)}{\mu} \frac{\partial p}{\partial r}$$
127

#2

$$k(\sigma) = \frac{1}{1 + \beta \sigma (1 + \beta_2 \sigma / \beta)}$$
129

Different empirical formulas for permeability decrease with an 130 increase in retained concentration have been proposed in the lit- 131 erature [14,22–26]. Usually, a linear function of  $\sigma$  in the denomi- 132 nator of Eq. (4) with  $\beta_2=0$  is used in the expression for formation 133 damage function  $k(\sigma)$ , where  $\beta$  is called the formation damage 134 coefficient [9,10].

A more general form of the denominator in Eq. (4) with non- 136 zero  $\beta_2$  is used in Ref. [24] to adjust coreflood injectivity impair- 137 ment data. The case of Eq. (4) corresponds to the linear function 138 of formation damage coefficient  $\beta$  versus retention concentration 139  $\sigma$ , which corresponds to a quadratic polynomial for reciprocal to 140 formation damage function  $k(\sigma)$ . Further in the text,  $\beta_2$  is called 141 the second formation damage coefficient. 142

The coefficients  $\beta$  and  $\beta_2$  are empirical constants characterizing 143 formation damage; they depend on the rock and the injected sus- 144 pension properties. 145

Systems (1) and (2) consist of two equations for two 146 unknowns—c and  $\sigma$ . For the injection of constant concentration 147 suspension into a clean bed, systems (1) and (3) are subject to the 148 following initial and boundary conditions: 149

$$t = 0: c = \sigma = 0, \quad r = r_w: c = c^0$$
 (5) **150**

The boundary condition (5) is set at the wellbore  $r=r_w$  (Fig. 2). 151 Introduce dimensionless coordinate and time 152

$$X = \frac{r^2}{R_c^2}, \quad t_D = \frac{qt}{\pi\phi h R_c^2} \tag{6}$$

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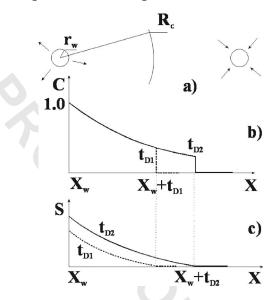


Fig. 2 Schematic for injected suspension propagation between injection and production well: (a) contour radius is equal to half-distance between injector and producer; (b) profile of suspension concentration is steady state behind the concentration front; (c) gradual accumulation of retained particles behind the concentration front

154 The dimensionless time with the pore volumes injected (p.v.i.) 155 unit is identical to that in waterflooding, while the linear coordi-**156** nate X is equal to the square of the dimensionless radius [26]. The **157** drainage radius is equal to the half-distance between injection and production wells (Fig. 2). 158

Substitution of Eq. (6) along with dimensionless concentra-159 160 tions, pressure, and filtration function

161 
$$C = \frac{c}{c^0}, \quad S = \frac{\sigma}{\phi c^0}, \quad p_D = \frac{4\pi h k_0 p}{q\mu}, \quad \Lambda(S) = \lambda(\sigma) R_c \quad (7)$$

**162** into systems (1), (2), and (4) yields

163

164

$$\frac{\partial C}{\partial t_D} + \frac{\partial C}{\partial X} = -\frac{\partial S}{\partial t_D}$$

 $\frac{\partial S}{\partial t_D} = \frac{\Lambda(S)}{2\sqrt{X}}C$ (9)

(8)

165 
$$1 = -\frac{X}{(1 + \beta\phi c^0 S + \beta_2(\phi c^0)^2 S^2)} \frac{\partial p_D}{\partial X}$$
(10)

Initial and boundary conditions (5) for dimensionless variables 166 167 (Eqs. (6) and (7)) take the form

**168** 
$$t_D = 0:C = S = 0$$
 and  $X = X_W = (r_W/R_c)^2:C = 1$  (11)

169 Several analytical solutions of the problem (Eqs. (1), (2), and **170** (5)) have been reported in the literature [6]. A semi-analytical 171 solution for the clean bed injection problem (Eq. (5)), for the case 172 of filtration and formation damage coefficients linear with respect 173 to retention concentration, is presented in the next section. 174 While the initial-boundary problem (Eqs. (1), (2), and (11)) is

**175** solved, pressure distribution along linear coordinate X during 176 flooding can be found from Eq. (4) by a direct integration of the 177 pressure gradient in X from the well radius  $X_w$  to the contour **178** radius  $X_c = 1$ .

#### 3 Semi-Analytical Model for Axisymmetric Deep Bed 179 Filtration 180

Let us derive a semi-analytical solution for the axisymmetric 181 (radial) problem of deep bed filtration for any arbitrary filtration 182 function  $\lambda(\sigma)$ . 183 184

Introduce potential

$$\Phi(S) = \int_0^S \frac{ds}{\Lambda(s)}$$
(12)

From Eq. (9), it follows that

$$C = \frac{\partial [2\sqrt{X\Phi(S)}]}{\partial t_D} \tag{13}$$

185

186

188

Substituting Eq. (13) in Eq. (8) yields

$$\frac{\partial}{\partial t_D} \left\{ \frac{\partial}{\partial t_D} [2\sqrt{X}\Phi(S)] \right\} + \frac{\partial}{\partial X} \left\{ \frac{\partial}{\partial t_D} [2\sqrt{X}\Phi(S)] \right\} = -\frac{\partial S}{\partial t_D} \quad (14)$$
189

Changing the order of differentiation in the second term of Eq. 190 (14) results in 191

$$\frac{\partial}{\partial t_D} \left\{ \frac{\partial}{\partial t_D} [2\sqrt{X}\Phi(S)] + \frac{\partial}{\partial X} [2\sqrt{X}\Phi(S)] + S \right\} = 0$$
 (15) 192

The integration of Eq. (15) in  $t_D$  accounting for initial condi- 193 tions (11) yields 194

$$\frac{\partial}{\partial t_D} [2\sqrt{X}\Phi(S)] + \frac{\partial}{\partial X} [2\sqrt{X}\Phi(S)] + S = \begin{cases} \frac{\partial}{\partial t_D} [2\sqrt{X}\Phi(S)] \\ \frac{\partial}{\partial t_D} [2\sqrt{X}\Phi(S)] \end{cases}$$
195

$$+ \frac{\partial}{\partial X} [2\sqrt{X}\Phi(S)] + S \bigg\}_{t_D=0}$$
(16) (16)

As it follows from initial conditions (11), the right hand side of 197 Eq. (16) is equal to zero; i.e., Eq. (16) becomes 198

$$\frac{\partial}{\partial t_D} [2\sqrt{X}\Phi(S)] + \frac{\partial}{\partial X} [2\sqrt{X}\Phi(S)] + S = 0$$
<sup>(17)</sup>
<sup>199</sup>

So, the introduction of potential (12) results in decreasing of the 200 order of the partial differential equation of mass balance (8) by 1. 201

Performing differentiation in Eq. (17) accounting for Eq. (12) 202 results in 203

$$\frac{\partial S}{\partial t_D} + \frac{\partial S}{\partial X} = \left(-\frac{\Phi(S)}{\sqrt{X}} - S\right) \frac{\Lambda(S)}{2\sqrt{X}}$$
(18)  
204

The obtained first order hyperbolic Eq. (18) can be solved by 205 method of characteristics. 206

Let us express first order partial differential Eq. (18) in charac- 207 teristic form, 208

$$\frac{dS}{dX} = \left(-\frac{\Phi(S)}{\sqrt{X}} - S\right) \frac{\Lambda(S)}{2\sqrt{X}}$$
(20)

From the boundary condition on the well (Eq. (11)), it follows 211 that for any moment  $t'_D$ 212

$$2\sqrt{X_w}\Phi(S) = t'_D \tag{21} 213$$

allowing expression of retained saturation on the wellbore 214

$$S(X_w, t'_D) = \Phi^{-1} \left( \frac{t'_D}{2\sqrt{X_w}} \right)$$
(22)  
215

Here,  $\Phi^{-1}$  is an inverse function to  $\Phi(S)$  (see Eq. (12)). 216

As it follows from Eq. (19), the characteristic line crossing any 217 arbitrary point  $(X, t_D)$  also crosses point  $(X_w, t'_D)$  (see also Fig. 3), **218** 

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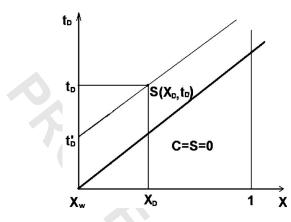


Fig. 3 Concentration front and characteristic line on the plane  $(X, t_D)$ 

**219** 
$$t'_D = t_D - X + X_w$$
 (23)

**220** It results in the Cauchy condition for the ordinary differential Eq. **221** (20),

$$X = X_{w}: S = \Phi^{-1} \left( \frac{t_{D} - X + X_{w}}{2\sqrt{X_{w}}} \right)$$
(24)

**223** Ordinary differential Eq. (20) subject to Cauchy condition (24) **224** can be solved numerically.

 Ahead of the concentration front  $X > X_f = t_D - X_w$ , suspended and retained concentrations are equal to zero. The retained con- centration is continuous along the front, S=0, while the suspen- sion concentration suffers discontinuity. Substituting Eq. (9) into Eq. (8) and accounting for zero retained concentration on the sus-pended concentration shock, we obtain

$$\frac{\partial C}{\partial t_D} + \frac{\partial C}{\partial X} = -\frac{\Lambda(0)}{2\sqrt{X}}C$$
(25)

**232** Front  $X=X_w+t_D$  is a characteristic line, which allows us to **233** calculate suspended concentration behind the front from Eq. (25),

234 
$$C(X_f(t_D), t_D) = \exp[-\Lambda(0)(\sqrt{X_f} - \sqrt{X_w})] = \exp[-\Lambda(0)(\sqrt{X_w} + t_D)]$$
  
235  $-\sqrt{X_w}$ ] (26)

**236** Consider the particular case of linear filtration coefficient:

$$\Lambda(S) = \lambda_0 R_c \left( 1 - \frac{S}{S_m} \right) \tag{27}$$

**238** The potential (12) is

222

231

237

241

$$\Phi(S) = -\frac{S_m}{\lambda_0 R_c} \ln\left(1 - \frac{S}{S_m}\right)$$
(2)

**240** Equation (20) and the Cauchy condition become

$$\frac{dS}{dX} = \left(\frac{S_m}{\lambda_0 R_c \sqrt{X}} \ln\left(1 - \frac{S}{S_m}\right) - S\right) \frac{\lambda_0 R_c}{2\sqrt{X}} \left(1 - \frac{S}{S_m}\right)$$
(29)

242 
$$X = X_w: S = S_m \left[ 1 - \exp\left(-\frac{\lambda_0 R_c}{S_m} \frac{t_D - X + X_w}{2\sqrt{X_w}}\right) \right]$$
(30)

 The structure of the suspension flow zone is shown in Fig. 3. Suspended and retained concentrations are equal to zero ahead of the front, which moves with carrier water velocity. Suspended concentration jumps from zero to the value of  $c^{-}(r_{f}, t)$  on the front. The suspended concentration behind the front is found from the condition on the characteristic line with  $\sigma=0$ , which coincides with the concentration front [26]:

250 
$$c^{-}(r_{f}(t),t) = c^{0} \exp\left[-\lambda(0)\left(\sqrt{r_{w}^{2} + \frac{qt}{\pi\phi h}} - r_{w}\right)\right]$$
 (31)

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The retained concentration behind the front is given by the **251** solution of the first order ordinary differential Eq. (20). For the **252** case of the linear filtration function (Eq. (27)), the solution is **253** given by formula (29). **254** 

The solution  $S(X, t_D)$  of the problem (Eqs. (29) and (30)) in the **255** behind-the-front region is obtained by the Runge–Kutta method. **256** The suspended concentration c(r,t) is obtained by Eq. (9) using **257** the retained concentration s(r,t). **258** 

Figure 4 presents profiles of suspended and retained concentra- 259 tions, (a) and (b), respectively, at five different moments. The 260 values of injectivity damage coefficients and other parameters are 261 obtained in Ref. [12] by the treatment of coreflood data [27]:  $\lambda_0$  262 = 10 m<sup>-1</sup>,  $\beta$ =100,  $\beta_2$ =-1000,  $\sigma_m$ =0.025,  $c_0$ =10 ppm,  $R_c$  263 = 500 m,  $r_w$ =0.1 m, and h=30 m. The axes X for the profiles 264 extend from well  $X=X_w$  until the value that corresponds to r 265 = 0.5 m. 266

The suspended concentration gradually increases with time in 267 each reservoir point for the calculated case of linear filtration co- 268 efficient. This behavior is typical of the declining filtration 269 coefficient—the outlet concentration after the breakthrough re- 270 mains constant for a constant filtration coefficient. Gradual accu- 271 mulation of retained concentration with time is shown in Fig. 4(*b*). 272 The dimensionless pressure drop between injector and reservoir 273

(drawdown) is 274

$$J = \frac{\Delta P}{\Delta P_0} = -\frac{1}{\ln(X_w)} \int_{X_w}^1 -\frac{\partial p_D}{\partial X} dX$$
(32)

282

Let us express pressure gradient via deposited concentration **276** from Eq. (10) and substitute it into Eq. (32), **277** 

$$J(t_D) = 1 - \frac{\beta \phi c^0}{\ln(X_w)} \int_{X_w}^1 \frac{S(X, t_D)}{X} dX - \frac{\beta_2(\phi c^0)^2}{\ln(X_w)} \int_{X_w}^1 \frac{S^2(X, t_D)}{X} dX$$
(33) 278

The well impedance (Eq. (33)) is obtained by numerical inte- **279** gration in *X* of expressions containing retained concentration dis- **280** tribution *S*(*X*, *t*<sub>D</sub>), as calculated from Eqs. (29) and (30). **281** 

### 4 Calculations of Well Injectivity Index

Let us show how to use formula (33) and the values of four **283** injectivity damage parameters  $\lambda_0$ ,  $\sigma_m$ ,  $\beta$ , and  $\beta_2$ , as obtained by **284** the treatment of coreflood data [12], for well injectivity decline **285** prediction. **286** 

The model for vertical well injectivity decline consists of non- **287** linear deep bed filtration equations for axisymmetric flow geom- **288** etry; see Eqs. (1), (3), and (4). An exact semi-analytical solution **289** of the radial deep bed filtration problem (Eqs. (29) and (30)) al- **290** lows us to calculate well impedance versus time (Eq. (33)). **291** 

Figure 5 shows the sensitivity of the impedance curve with 292 respect to the second formation damage coefficient, while the fil- 293 tration coefficient is constant; i.e.,  $\sigma_m$  tends to infinity. A 1 month 294 injection period in Fig. 5(*a*) corresponds to  $6 \times 10^{-3}$  p.v.i. in Fig. 295 5(*b*). The default case corresponds to linear filtration with constant 296 filtration and formation damage coefficients,  $\lambda_0=10 \text{ m}^{-1}$ ,  $\beta=50$ , 297 and  $\beta_2=0$ . Curves 2 and 3 correspond to values of the second 298 formation damage coefficient  $\beta_2=200$  and  $\beta_2=-200$ , respectively. 299 The  $\beta_2$ -values are in the range of those obtained from coreflood 300 data [27]. 301

If compared with the default case with  $\beta_2=0$ , adding the qua- 302 dratic term  $\beta_2\sigma^2$  with the positive  $\beta_2$  coefficient results in imped- 303 ance increase (curve 2). Introduction of the quadratic term  $\beta_2\sigma^2$  304 with the negative  $\beta_2$  coefficient results in impedance decrease 305 (curve 3). The impedance curve is concave for the case of the 306 positive second formation damage coefficient and is convex for 307 negative  $\beta_2$ . The significant difference between the impedance 308

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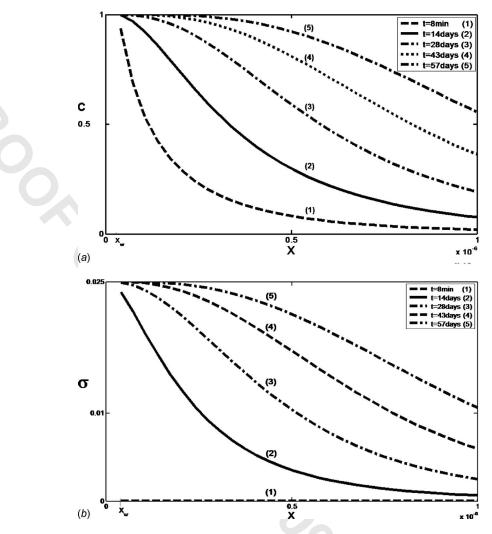


Fig. 4 Dynamics of (a) suspended and (b) retained concentration profiles during suspension injection in vertical well

**309** curves with zero, positive, and negative  $\beta_2$  starts to appear after **310**  $2 \times 10^{-3}$  p.v.i. (10 days). The curves almost coincide before the **311** injection of  $10^{-3}$  p.v.

**312** After 1 month of injection, the injectivity index decreases 2.4 **313** times for the constant formation damage coefficient ( $\beta_2=0$ ). It

**314** decreases 2.65 times for  $\beta_2 = 200$  and 2.15 times for  $\beta_2 = -200$ .

 Let us discuss the effect of the nonconstant filtration coefficient on impedance growth. Figure 6 presents impedance curves for constant formation damage coefficients  $\beta$ =50 and  $\beta_2$ =0. The ini- tial filtration coefficient  $\lambda_0$ =10 m<sup>-1</sup>. Curves 1–5 correspond to the following values for maximum retained concentration: infinity, 0.1, 0.05, 0.025, and 0.0015, respectively. So, the filtration func- tion decreases in order of the curve number increase. The larger the filtration function, the higher the retained particle concentra- tion and the higher the impedance. Therefore, curves 1–5 follow in the order of the impedance decrease.

 Curve 1 is linear due to the constant filtration coefficient. Im- pedance curves 2–5 stabilize with time; i.e., its time derivative tends to zero when time tends to infinity. It is explained by reach- ing the maximum retention concentration  $\sigma_m$  in each point around the well. The higher the distance to the point from the injector, the later the maximum retention concentration will be reached. There- fore, curves 2–5 follow in order of decrease in the stabilization time. The moments of impedance stabilization are 200 days, 106 days, 50 days, and 3 days, for curves 2–5, respectively.

334 Since the impedance function asymptotically reaches maximum

for retained concentration that is equal to the critical value  $\sigma_m$ , 335 from Eq. (33) follows formula for the maximum impedance, 336

$$J(t \to \infty) = 1 + \beta \sigma_m + \beta_2 \sigma_m^2 \tag{34}$$

After the long-term injection, the injectivity index decreases **338** 6.0, 3.5, 2.3, and 1.1 times for curves 2–5, respectively. **339** 

In offshore waterflood operations, the injectivity index stabili- 340 zation at some reduced value happens quite often [3,9,19]. The 341 usual explanation is well fracturing, where the fracture opens due 342 to high pressure near the injector because of reduced permeability; 343 it propagates into the formation during further permeability de-344 cline and increase in hydraulic resistivity for leak-off. Another 345 explanation is erosion of external and internal filter cakes, where 346 the pressure gradient, increased due to permeability decline, drags 347 particles from cakes in situ porous media and on the inner well 348 surface. 349

Reaching the maximum retention concentration, where particle **350** capture by the rock does not happen anymore, is another explana-**351** tion of well injectivity index stabilization. **352** 

Let us compare the impedance growth curves for constant and **353** for linear filtration and formation damage functions. Straight line **354** 1 in Fig. 7 corresponds to constant filtration and formation dam-**355** age coefficients:  $\lambda_0 = 10 \text{ m}^{-1}$  and  $\beta = 50$ . For curves 2–6, values of **356**  $\lambda_0$  and  $\beta$  are the same as that for case 1.

The second formation damage coefficient for curve 2 is equal to **358** 1000. For small retention concentrations at the beginning of in- **359** 

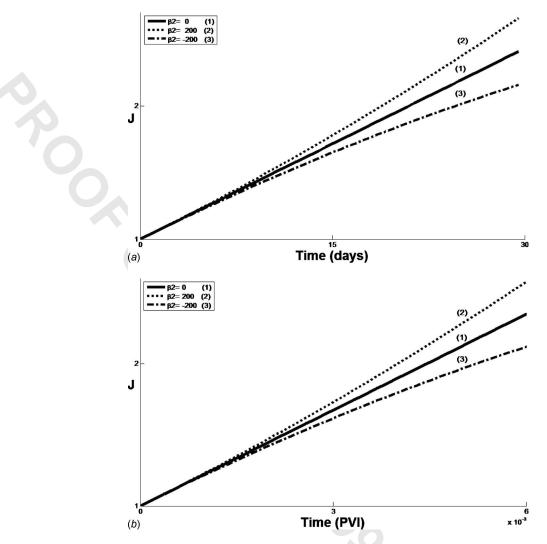


Fig. 5 Effect of two formation damage coefficients  $\beta$  and  $\beta_2$  on well impedance curve: (*a*) impedance versus real time; (*b*) impedance versus p.v.i.

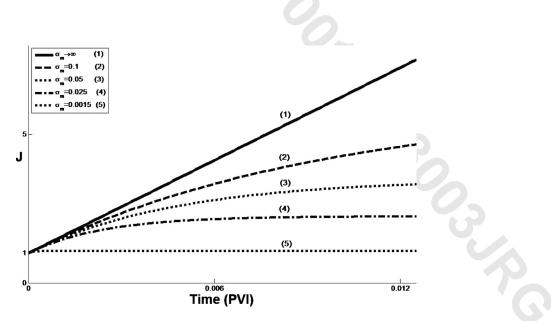


Fig. 6 Sensitivity of impedance curve to variation in filtration function: impedance versus p.v.i.

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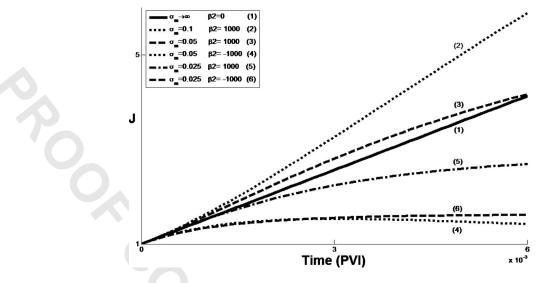


Fig. 7 Monotonic and nonmonotonic impedance curves

 jection, the values of filtration functions for cases 1 and 2 are almost equal. Due to the high value of the second formation dam- age coefficient for curve 2, it is located above the straight line 1 for small times. Curve 2 stabilizes with time, while the straight line 1 grows unlimitedly. Therefore, at some moment, curve 2 intersects line 1 and tends to a constant value at large times.

 Curve 3 has a lower maximum retention concentration value if compared with curve 2. Therefore, curve 3 is located above line 1 and below curve 2 at small times. Curve 3 crosses line 1 and stabilize faster than curve 2.

The value of maximum retention concentration for curve 5 is is lower than that for curve 3. Curve 5 is located above line 1 and is below the curve 3 at small times. Curve 5 crosses line 1 and stabilize faster than curve 3.

**374** Curves 2, 3, and 5 exhibit a monotonic growth. For positive **375**  $\sigma$ -values and positive second formation damage coefficient, the **376** derivative of reciprocal to formation damage function  $k(\sigma)$  (see **377** Eq. (4))

$$(\beta_2 \sigma^2 + \beta \sigma + 1)' = 2\beta_2 \sigma + \beta > 0$$

**379** is always positive; i.e., the hydraulic resistance of retained par-**380** ticles increases during the deposition.

 Curves 4 and 6 correspond to the negative second formation damage coefficient  $\beta_2 = -1000$ . The derivative of the formation damage function  $2\beta_2\sigma + \beta$  is positive for low  $\sigma$ -values, i.e., for  $\sigma < -\beta/2\beta_2$  (Fig. 8). For higher retention concentration values,

> $k_r^{-1}(\sigma)$ 1.0 0  $-\beta/2\beta_2$   $\sigma_m$   $\sigma$

Fig. 8 Nonmonotonic behavior of reciprocal to formation damage function due to linear interpolation of formation damage coefficient

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 $\sigma > -\beta/2\beta_2$ , the derivative is negative; accumulation of retained **385** particles results in the reduction of hydraulic resistivity. This effect is physically unrealistic. **387** 

In case 4,  $\sigma_m$ =0.05,  $\beta_2$ =-1000, and  $\beta$ =50; so,  $-\beta/2\beta_2$  388 =0.025, which is less than  $\sigma_m$ ; i.e., the hydraulic resistivity in-389 creases for retention concentration varying from zero to 0.025 and 390 decreases afterward for retention concentration varying from 391 0.025 to 0.05. Curve 4 decreases for  $\sigma$ >0.025. This behavior is 392 unrealistic. The paradox is caused by approximation of the forma-393 tion damage function by the quadratic polynomial based on 394 *J*-values for small retention concentrations (Fig. 8). If the depen-395 dency  $\lambda = \lambda(\sigma)$  has a negative second derivative, its approximation 396 by the linear function (3) causes the high value for the maximum 397 retained concentration (Fig. 9), resulting in inequality  $\sigma_m$ > 398  $-\beta/2\beta_2$ . In this case, the model with linear filtration coefficient 399 (3) is not valid.

For case 6,  $\sigma_m = 0.025$  and  $-\beta/2\beta_2 = 0.025 = \sigma_m$ . During an in- 401 crease in retention concentration from zero to  $\sigma_m$ , where  $\sigma$  tends 402 to  $\sigma_m$  asymptotically, the hydraulic resistivity increases, which is 403 shown by the behavior of curve 6. 404

Figure 10 presents injectivity index decline for the cases dis- 405 cussed in Fig. 7. The declined form of injectivity index versus 406 time does allow distinguishing linear and nonlinear well behavior, 407 while the impedance plot clearly shows a linear well behavior for 408 the case of constant injectivity damage coefficients and nonlinear 409

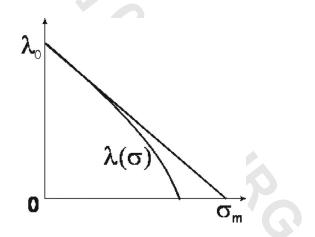


Fig. 9 Overestimated value of maximum retained concentration due to linear approximation of the filtration function

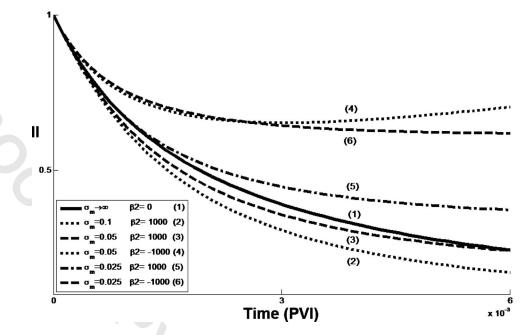


Fig. 10 Injectivity decline analysis using well index curves

**410** curves for varying  $\lambda$  and  $\beta$ . It shows the advantage to analyze 411 injectivity impairment in coordinates "impedance versus time." 412 The proposed three-point-pressure method [12] can be imple-413 mented into a simple, robust, and compact tool for applications in 414 on-site field conditions. Figure 11 shows the use of the tool at the 415 sea platform. Direct measurements data of the rate and of pres-**416** sures in three core points are treated by the computer program **417** with optimization algorithm minimizing the deviation between the 418 modeling and coreflood data. It allows us to determine the four **419** injectivity damage coefficients  $\lambda_0$ ,  $\sigma_m$ ,  $\beta$ , and  $\beta_2$  from coreflood 420 data. Then, the impedance and well index are recalculated for the 421 axisymmetric flow of a vertical well (Eq. (33)) and also for flow geometries of horizontal, fractured, and perforated wells using the 422 analytical solutions for 3D flow problems [28]. 423

### **424 5** Conclusions

425 Injectivity decline due to rock clogging by the injected suspension with linear filtration and formation damage functions allows 426 for semi-analytical modeling. 427

428 The impedance for injectivity damage parameters that are linear 429 functions of the retained concentration monotonically increases **430** and asymptotically tends to a limit constant, while the impedance 431 for constant injectivity damage parameters linearly grows with 432 time.

The injectivity index stabilizes after the retained particle con-433 434 centration reaches its maximum value in some neighborhood of 435 the injector.



Fig. 11 Using the three-point-pressure tool at sea platform

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The well behavior for constant injectivity damage parameters 436 can be distinguished from that for linear function parameters by 437 increasing impedance curves, not by declining well index curves. 438

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## Acknowledgment

The authors thank Petrobras for generous sponsorship of the 440 research project during many years and Dr. Farid Shecaira and 441 Eng. Alexandre Guedes Siqueira (Petrobras/CENPES) for fruitful 442 support and encouragement. Many thanks are due to Prof. P. Cur- 443 rie (Delft University of Technology), Prof. Yannis Yortsos (Uni- 444 versity of Southern California), and Prof. A. Shapiro for fruitful 445 discussions. 446

Nomenclature		447
Latin Letters		448
A =	specific rock surface, $[L]^2$	<b>450</b>
b =	area on grain surface filled by one retained	451
	particle, $[L]^2$	452
<i>c</i> =	suspension particle concentration	453
	injected suspension concentration	454
h =	formation thickness, [L]	456
II =	injectivity index	457
	impedance	458
k =	permeability, $[L]^2$	469
	core length, [L]	462
p =	pressure, $[M][T]^{-2}[L]^{-1}$	463
r =	radius, [L]	465
$R_c =$	drainage (contour) radius, [L]	468
q =	injection rate, $[L^3/T]$	460
t =	time, $[T]$	472
U =	Darcy velocity, $[L][T]^{-1}$	478
	coordinate in linear geometry, $[L]$	476
	dimensionless coordinate in radial geometry	477

## **Greek Letters**

- $\beta$  = formation damage coefficient  $\beta_2$  = formation damage coefficient
- λ = filtration coefficient  $[L]^{-1}$
- $\lambda_0$  = value of filtration coefficient for  $\sigma=0, [L]$ 483

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485	$\mu$ = viscosity of water, $[M][T]^{-1}[L]^{-1}$
487	$\sigma$ = deposited particle concentration
488	$\Phi = \text{porosity}$
489	$\Phi$ = potential
490 Supe	erscripts/Subscripts
491	0 = initial
102	f = front

ess

- = well w

#### **496** References

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